# Wavelet Transform Associated with the $q$-Dunkl Operator 

Néji Bettaibi, Rym H. Bettaieb ${ }^{\dagger}$<br>Institut Préparatoire aux Études d'Ingénieur de Monastir, 5000 Monastir, Tunisia<br>and<br>Slim Bouaziz ${ }^{\ddagger}$<br>Institut Préparatoire aux Études d'Ingénieur Elmanar, 1060 Tunis, Tunisia

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#### Abstract

In this paper, we present some new elements of harmonic analysis related to the $q$-Dunkl operator introduced in [1], we define and study the $q$-wavelets and the continuous $q$-wavelet transforms associated with this operator. Next, as an application, we give inversion formulas for the $q$-Dunkl intertwining operator and its dual using $q$-wavelets.


## 1. Introduction

In [11, 12], R. L. Rubin constructed a $q^{2}$-analogue Fourier analysis associated with a $q^{2}$-analogue differential operator $\partial_{q}$. Using this $q$-harmonic analysis, the

[^0]authors studied in [4] the $q$-wavelets and the continuous $q$-wavelet transforms associated with the operator $\partial_{q}$.

In [1], the authors introduced a $q$-analogue of the Dunkl operator on $\mathbb{R}$ and they defined and studied its associated Fourier transform $F_{D}^{\alpha, q}$, called $q$-Dunkl transform, which is a $q$-analogue of the Bessel-Dunkl transform. They, also, studied the $q$-Dunkl intertwining operator $V_{\alpha, q}$ and its dual ${ }^{t} V_{\alpha, q}$ via the $q$-analogues of the Riemann-Liouville and Weyl transforms $R_{\alpha, q}$ and ${ }^{t} R_{\alpha, q}$, studied in [5]. In particular, they proved that $V_{\alpha, q}$ and its dual are automorphism of some spaces $\mathcal{E}_{q}\left(\mathbb{R}_{q}\right)$ and $\mathcal{D}_{q}\left(\mathbb{R}_{q}\right)$, respectively and they gave their inversion operators using $R_{\alpha, q}$ and ${ }^{t} R_{\alpha, q}$.

In this paper, we define the generalized $q$-Dunkl translation operator and its related convolution product, we give some of their properties, then, we are interested by studying the $q$-wavelets and the continuous $q$-wavelet transforms associated with the $q$-Dunkl operator. Next, we establish an inversion formulas for the $q$-Dunkl intertwining operator $V_{\alpha, q}$ and its dual ${ }^{t} V_{\alpha, q}$ using $q$-wavelets.

This paper is organized as follows: in Section 2, we present some preliminaries results and notations that will be useful in the sequel. In Section 3, we recall some results and properties concerning the $q$-Dunkl transform studied in [1], we introduce the generalized $q$-Dunkl translation operator and its related convolution product and we give some of their properties. In Section 4 , we define and study the $q$-wavelet and the continuous $q$-wavelet transform associated with the $q$-Dunkl operator, and we provide for this transform a Plancherel formula and an inversion theorem. Section 5 is devoted to give some inversion formulas for the $q$-Dunkl intertwining operator and its dual on some new spaces (other than $\mathcal{E}_{q}\left(\mathbb{R}_{q}\right)$ and $\mathcal{D}_{q}\left(\mathbb{R}_{q}\right)$ ). Finally, in Section 6, we give some relations between the continuous $q$-wavelet transform associated with the $q$-Dunkl operator and those associated with the $q^{2}$-analogue differential operator $\partial_{q}$, studied in [4]. Next, by the help of these relations, we derive the inversion formulas of the $q$-Dunkl intertwining operator and its dual using $q$-wavelets.

## 2. Notations and Preliminaries

For the convenience of the reader, we provide in this section a summary of the mathematical notations and definitions used in this paper, we will follow the notations of $[11,12]$. We fix $q \in] 0,1[$ and we refer to the book by G. Gasper and M. Rahman [6] for the definitions, notations and properties of the $q$-shifted factorials and the $q$-hypergeometric functions.

We will write
$\mathbb{R}_{q}=\left\{ \pm q^{n} \quad: \quad n \in \mathbb{Z}\right\}, \widetilde{\mathbb{R}}_{q}=\left\{ \pm q^{n} \quad: \quad n \in \mathbb{Z}\right\} \cup\{0\}$.
For $a \in \mathbb{C}$, the $q$-shifted factorials are defined by

$$
\begin{equation*}
(a ; q)_{0}=1 ; \quad(a ; q)_{n}=\prod_{k=0}^{n-1}\left(1-a q^{k}\right), \quad n=1,2, \ldots ; \quad(a ; q)_{\infty}=\prod_{k=0}^{\infty}\left(1-a q^{k}\right) \tag{1}
\end{equation*}
$$

We also denote

$$
\begin{equation*}
[x]_{q}=\frac{1-q^{x}}{1-q}, \quad x \in \mathbb{C} ; \quad[n]_{q}!=\frac{(q ; q)_{n}}{(1-q)^{n}}, \quad n \in \mathbb{N} . \tag{2}
\end{equation*}
$$

The $q^{2}$-analogue differential operator is (see [12], [11]),

$$
\partial_{q}(f)(z)=\left\{\begin{array}{cl}
\frac{f\left(q^{-1} z\right)+f\left(-q^{-1} z\right)-f(q z)+f(-q z)-2 f(-z)}{2(1-q) z} & \text { if } z \neq 0  \tag{3}\\
\lim _{x \rightarrow 0} \partial_{q}(f)(x) \quad\left(\text { in } \mathbb{R}_{q}\right) & \text { if } z=0
\end{array}\right.
$$

Note that if $f$ is differentiable at $z$, then $\lim _{q \rightarrow 1} \partial_{q}(f)(z)=f^{\prime}(z)$.
The $q$-trigonometric functions $q$-cosine and $q$-sine are defined by (see $[11,12]$ ):

$$
\begin{equation*}
\cos \left(x ; q^{2}\right)=\sum_{n=0}^{\infty}(-1)^{n} q^{n(n+1)} \frac{x^{2 n}}{[2 n]_{q}!} \tag{4}
\end{equation*}
$$

and

$$
\begin{equation*}
\sin \left(x ; q^{2}\right)=\sum_{n=0}^{\infty}(-1)^{n} q^{n(n+1)} \frac{x^{2 n+1}}{[2 n+1]_{q}!} . \tag{5}
\end{equation*}
$$

These functions induce a $\partial_{q^{-}}$-adapted $q^{2}$-analogue exponential function by

$$
\begin{equation*}
e\left(z ; q^{2}\right)=\cos \left(-i z ; q^{2}\right)+i \sin \left(-i z ; q^{2}\right) \tag{6}
\end{equation*}
$$

$e\left(z ; q^{2}\right)$ is absolutely convergent for all $z$ in the plane since both of its component functions are. $\lim _{q \rightarrow 1^{-}} e\left(z ; q^{2}\right)=e^{z}$ (exponential function) pointwise and uniformly on compacts.
Using the same technique as in [11], one can prove that for all $x \in \mathbb{R}_{q}$, we have

$$
\left|\cos \left(x ; q^{2}\right)\right| \leq \frac{1}{(q ; q)_{\infty}} \quad \text { and } \quad\left|\sin \left(x ; q^{2}\right)\right| \leq \frac{1}{(q ; q)_{\infty}}
$$

so,

$$
\begin{equation*}
\forall x \in \mathbb{R}_{q}, \quad\left|e\left(-i x ; q^{2}\right)\right| \leq \frac{2}{(q ; q)_{\infty}} \tag{7}
\end{equation*}
$$

The $q$-Jackson integrals from 0 to $a \in \mathbb{R}$ and from $-\infty$ to $+\infty$ are defined by (see [7], [8], [10], [9])

$$
\begin{equation*}
\int_{0}^{a} f(x) d_{q} x=(1-q) a \sum_{n=0}^{\infty} f\left(a q^{n}\right) q^{n} \tag{8}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{-\infty}^{+\infty} f(x) d_{q} x=(1-q) \sum_{n=-\infty}^{\infty}\left\{f\left(q^{n}\right)+f\left(-q^{n}\right)\right\} q^{n} \tag{9}
\end{equation*}
$$

provided the sums converge absolutely.
The following result can be verified by direct computation.
Lemma 1. If $\int_{-\infty}^{\infty} f(t) d_{q} t$ exists, then for all $a \in \mathbb{R}_{q}$,

$$
\int_{-\infty}^{\infty} f(a t) d_{q} t=|a|^{-1} \int_{-\infty}^{\infty} f(t) d_{q} t
$$

In the sequel, we will need the following spaces:

- $\mathcal{E}_{q}\left(\mathbb{R}_{q}\right)$ the space of functions $f$ defined on $\mathbb{R}_{q}$, satisfying
$\forall n \in \mathbb{N}, \quad a \geq 0, \quad P_{n, a}(f)=\sup \left\{\left|\partial_{q}^{k} f(x)\right| ; 0 \leq k \leq n ; x \in[-a, a] \cap \mathbb{R}_{q}\right\}<\infty$ and

$$
\lim _{x \rightarrow 0} \partial_{q}^{n} f(x) \quad\left(\text { in } \quad \mathbb{R}_{q}\right) \quad \text { exists. }
$$

We provide it with the topology defined by the semi norms $P_{n, a}$.

- $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ the space of functions $f$ defined on $\mathbb{R}_{q}$ satisfying

$$
\forall n, m \in \mathbb{N}, \quad P_{n, m, q}(f)=\sup _{x \in \mathbb{R}_{q}}\left|x^{m} \partial_{q}^{n} f(x)\right|<+\infty
$$

and

$$
\lim _{x \rightarrow 0} \partial_{q}^{n} f(x) \quad\left(\text { in } \quad \mathbb{R}_{q}\right) \quad \text { exists. }
$$

- $\mathcal{D}_{q}\left(\mathbb{R}_{q}\right)$ the subspace of $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ constituted of functions with compact supports.
- $L_{\alpha, q}^{p}\left(\mathbb{R}_{q}\right)=\left\{f:\|f\|_{p, \alpha, q}=\left(\int_{-\infty}^{\infty}|f(x)|^{p}|x|^{2 \alpha+1} d_{q} x\right)^{\frac{1}{p}}<\infty\right\}, \quad p>0$ and $\alpha \in \mathbb{R}$.
- $L_{q}^{\infty}\left(\mathbb{R}_{q}\right)=\left\{f:\|f\|_{\infty, q}=\sup _{x \in \mathbb{R}_{q}}|f(x)|<\infty\right\}$.


## 3. Elements of $q$-Dunkl Harmonic Analysis

In this section, we collect some basic results and properties from the $q$-Dunkl operator theory, studied in [1], and we introduce and study a generalized $q$ Dunkl translation as well as its related convolution product.

For $\alpha \geq-\frac{1}{2}$, the $q$-Dunkl transform is defined on $L_{\alpha, q}^{1}\left(\mathbb{R}_{q}\right)$ by (see [1])

$$
\begin{equation*}
F_{D}^{\alpha, q}(f)(\lambda)=c_{\alpha, q} \int_{-\infty}^{+\infty} f(x) \psi_{-\lambda}^{\alpha, q}(x) \cdot|x|^{2 \alpha+1} d_{q} x, \quad \lambda \in \widetilde{\mathbb{R}}_{q} \tag{10}
\end{equation*}
$$

where $c_{\alpha, q}=\frac{(1+q)^{-\alpha}}{2 \Gamma_{q^{2}}(\alpha+1)} \quad$ and $\quad \psi_{\lambda}^{\alpha, q}$ is the $q$-Dunkl kernel defined by

$$
\begin{equation*}
\psi_{\lambda}^{\alpha, q}: x \longmapsto j_{\alpha}\left(\lambda x ; q^{2}\right)+\frac{i \lambda x}{[2 \alpha+2]_{q}} j_{\alpha+1}\left(\lambda x ; q^{2}\right), \tag{11}
\end{equation*}
$$

with $j_{\alpha}\left(x ; q^{2}\right)$ is the normalized third Jackson's $q$-Bessel function given by:
$j_{\alpha}\left(x ; q^{2}\right)=\sum_{n=0}^{+\infty}(-1)^{n} \frac{\Gamma_{q^{2}}(\alpha+1) q^{n(n+1)}}{\Gamma_{q^{2}}(\alpha+n+1) \Gamma_{q^{2}}(n+1)}\left(\frac{x}{1+q}\right)^{2 n}$.

It was proved in [1] that for all $\lambda \in \mathbb{C}$, the function: $\quad x \longmapsto \psi_{\lambda}^{\alpha, q}(x)$ is the unique solution of the $q$-differential-difference equation:

$$
\left\{\begin{array}{cl}
\Lambda_{\alpha, q}(f) & =i \lambda f  \tag{12}\\
f(0) & =1
\end{array}\right.
$$

where $\Lambda_{\alpha, q}$ is the $q$-Dunkl operator defined by

$$
\begin{equation*}
\Lambda_{\alpha, q}(f)(x)=\partial_{q}\left[f_{e}+q^{2 \alpha+1} f_{o}\right](x)+[2 \alpha+1]_{q} \frac{f(x)-f(-x)}{2 x}, \tag{13}
\end{equation*}
$$

with $f_{e}$ and $f_{o}$ are respectively the even and the odd parts of $f$.
We recall that the $q$-Dunkl operator $\Lambda_{\alpha, q}$ lives the spaces $\mathcal{D}_{q}\left(\mathbb{R}_{q}\right)$ and $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ invariant. Some other properties of the $q$-Dunkl kernel and the $q$-Dunkl transform are given in the following results (see [1]).

## Proposition 1.

i) $\psi_{\lambda}^{\alpha, q}(x)=\psi_{x}^{\alpha, q}(\lambda), \quad \psi_{a \lambda}^{\alpha, q}(x)=\psi_{\lambda}^{\alpha, q}(a x), \quad \overline{\psi_{\lambda}^{\alpha, q}(x)}=\psi_{-\lambda}^{\alpha, q}(x), \forall \lambda, x \in \mathbb{R}$, $a \in \mathbb{C}$.
ii) If $\alpha=-\frac{1}{2}$, then $\psi_{\lambda}^{\alpha, q}(x)=e\left(i \lambda x ; q^{2}\right)$.

For $\alpha>-\frac{1}{2}, \quad \psi_{\lambda}^{\alpha, q}$ has the following $q$-integral representation of Mehler type

$$
\begin{equation*}
\psi_{\lambda}^{\alpha, q}(x)=\frac{(1+q) \Gamma_{q^{2}}(\alpha+1)}{2 \Gamma_{q^{2}}\left(\frac{1}{2}\right) \Gamma_{q^{2}}\left(\alpha+\frac{1}{2}\right)} \int_{-1}^{1} \frac{\left(t^{2} q^{2} ; q^{2}\right)_{\infty}}{\left(t^{2} q^{2 \alpha+1} ; q^{2}\right)_{\infty}}(1+t) e\left(i \lambda x t ; q^{2}\right) d_{q} t \tag{14}
\end{equation*}
$$

iii) For all $\lambda \in \mathbb{R}_{q}, \psi_{\lambda}^{\alpha, q}$ is bounded on $\widetilde{\mathbb{R}}_{q}$ and we have

$$
\begin{equation*}
\left|\psi_{\lambda}^{\alpha, q}(x)\right| \leq \frac{4}{(q ; q)_{\infty}}, \quad \forall x \in \widetilde{\mathbb{R}}_{q} \tag{15}
\end{equation*}
$$

iv) For all $\lambda \in \mathbb{R}_{q}, \psi_{\lambda}^{\alpha, q} \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$.
v) The function $\psi_{\lambda}^{\alpha, q}$ verifies the following orthogonality relation: For all $x, y \in$ $\mathbb{R}_{q}$,

$$
\begin{equation*}
\int_{-\infty}^{+\infty} \psi_{\lambda}^{\alpha, q}(x) \overline{\psi_{\lambda}^{\alpha, q}(y)}|\lambda|^{2 \alpha+1} d q \lambda=\frac{4(1+q)^{2 \alpha} \Gamma_{q^{2}}^{2}(\alpha+1) \delta_{x, y}}{(1-q)|x y|^{\alpha+1}} \tag{16}
\end{equation*}
$$

vi) If $f \in L_{\alpha, q}^{1}\left(\mathbb{R}_{q}\right)$ then $F_{D}^{\alpha, q}(f) \in L_{q}^{\infty}\left(\mathbb{R}_{q}\right)$,

$$
\begin{equation*}
\left\|F_{D}^{\alpha, q}(f)\right\|_{\infty, q} \leq \frac{4 c_{\alpha, q}}{(q ; q)_{\infty}}\|f\|_{1, \alpha, q} \tag{17}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{\substack{|\lambda| \rightarrow+\infty \\ \lambda \in \mathbb{R}_{q}}} F_{D}^{\alpha, q}(f)(\lambda)=0, \quad \lim _{\substack{|\lambda| \rightarrow 0 \\ \lambda \in \mathbb{\mathbb { R }}_{q}}} F_{D}^{\alpha, q}(f)(\lambda)=F_{D}^{\alpha, q}(f)(0) \tag{18}
\end{equation*}
$$

vii) For $f \in L_{\alpha, q}^{1}\left(\mathbb{R}_{q}\right)$,

$$
\begin{equation*}
F_{D}^{\alpha, q}\left(\Lambda_{\alpha, q} f\right)(\lambda)=i \lambda F_{D}^{\alpha, q}(f)(\lambda) \tag{19}
\end{equation*}
$$

viii) For $f, g \in L_{\alpha, q}^{1}\left(\mathbb{R}_{q}\right)$,

$$
\begin{equation*}
\int_{-\infty}^{+\infty} F_{D}^{\alpha, q}(f)(\lambda) g(\lambda)|\lambda|^{2 \alpha+1} d_{q} \lambda=\int_{-\infty}^{+\infty} f(x) F_{D}^{\alpha, q}(g)(x)|x|^{2 \alpha+1} d_{q} x \tag{20}
\end{equation*}
$$

Theorem 1. For all $f \in L_{\alpha, q}^{1}\left(\mathbb{R}_{q}\right)$, we have

$$
\begin{align*}
\forall x \in \mathbb{R}_{q}, \quad f(x) & =c_{\alpha, q} \int_{-\infty}^{+\infty} F_{D}^{\alpha, q}(f)(\lambda) \psi_{\lambda}^{\alpha, q}(x) \cdot|\lambda|^{2 \alpha+1} d_{q} \lambda  \tag{21}\\
& =\overline{F_{D}^{\alpha, q}\left(\overline{F_{D}^{\alpha, q}(f)}\right)}(x)
\end{align*}
$$

Theorem 2. i) Plancherel formula
For $\alpha \geq-1 / 2$, the $q$-Dunkl transform $F_{D}^{\alpha, q}$ is an isomorphism from $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ onto itself. Moreover, for all $f \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, we have

$$
\begin{equation*}
\left\|F_{D}^{\alpha, q}(f)\right\|_{2, \alpha, q}=\|f\|_{2, \alpha, q} \tag{22}
\end{equation*}
$$

## ii) Plancherel theorem

The $q$-Dunkl transform can be uniquely extended to an isometric isomorphism on $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$. Its inverse transform $\left(F_{D}^{\alpha, q}\right)^{-1}$ is given by :

$$
\begin{equation*}
\left(F_{D}^{\alpha, q}\right)^{-1}(f)(x)=c_{\alpha, q} \int_{-\infty}^{+\infty} f(\lambda) \psi_{\lambda}^{\alpha, q}(x) \cdot|\lambda|^{2 \alpha+1} d_{q} \lambda=F_{D}^{\alpha, q}(f)(-x) \tag{23}
\end{equation*}
$$

We are now in a position to define the generalized $q$-Dunkl translation operator.

Definition 1. The generalized $q$-Dunkl translation operator is defined for $f \in$ $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$ and $x, y \in \mathbb{R}_{q}$ by

$$
\begin{align*}
T_{y}^{\alpha ; q}(f)(x) & =c_{\alpha, q} \int_{-\infty}^{\infty} F_{D}^{\alpha, q}(f)(\lambda) \psi_{\lambda}^{\alpha, q}(x) \psi_{\lambda}^{\alpha, q}(y)|\lambda|^{2 \alpha+1} d_{q} \lambda,  \tag{24}\\
T_{0}^{\alpha ; q}(f) & =f .
\end{align*}
$$

It verifies the following properties.

## Proposition 2.

1) For all $f \in L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$ and all $x \in \mathbb{R}_{q}$,

$$
\lim _{\substack{y \rightarrow 0 \\ y \in \mathbb{R}_{q}}} T_{y}^{\alpha ; q}(f)(x)=f(x) .
$$

2) For all $x, y \in \mathbb{R}_{q}, T_{y}^{\alpha ; q}(f)(x)=T_{x}^{\alpha ; q}(f)(y)$.
3) If $f \in L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right) \quad$ (resp. $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ ) then $T_{y}^{\alpha ; q}(f) \in L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right) \quad$ (resp. $\left.\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)\right)$ and we have

$$
\begin{equation*}
\left\|T_{y}^{\alpha ; q}(f)\right\|_{2, \alpha, q} \leq \frac{4}{(q ; q)_{\infty}}\|f\|_{2, \alpha, q} \tag{25}
\end{equation*}
$$

4) For all $x, y, \lambda \in \mathbb{R}_{q}, T_{y}^{\alpha ; q}\left(\psi_{\lambda}^{\alpha, q}\right)(x)=\psi_{\lambda}^{\alpha, q}(x) \psi_{\lambda}^{\alpha, q}(y)$.
5) For $f \in L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right), x, y \in \mathbb{R}_{q}$, we have

$$
\begin{equation*}
F_{D}^{\alpha, q}\left(T_{y}^{\alpha ; q} f\right)(\lambda)=\psi_{\lambda}^{\alpha, q}(y) F_{D}^{\alpha, q}(f)(\lambda) \tag{26}
\end{equation*}
$$

6) For $f \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ and $y \in \mathbb{R}_{q}$, we have

$$
\Lambda_{\alpha, q} T_{y}^{\alpha ; q} f=T_{y}^{\alpha ; q} \Lambda_{\alpha, q} f
$$

## Proof.

1) Since $\psi_{\lambda}^{\alpha, q}$ is bounded on $\mathbb{R}_{q}$, then the Lebesgue theorem and Theorem 1 give the result.
2) Follows from the definition of the generalized $q$-Dunkl translation.
3) Since $f \in L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right) \quad$ (resp. $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ ), $\psi_{y}^{\alpha, q}$ is bounded for all $y$ (resp. $\psi_{y}^{\alpha, q} \in$ $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ ) and $F_{D}^{\alpha, q}$ is an automorphism of $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right) \quad\left(\right.$ resp. $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ ), then $F_{D}^{\alpha, q}(f) \cdot \psi_{y}^{\alpha, q}$ is in $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right) \quad\left(\right.$ resp. $\left.\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)\right)$. So, by the help of the Plancherel theorem, we have for all $y \in \mathbb{R}_{q}, T_{y}^{\alpha ; q} f=\left(F_{D}^{\alpha, q}\right)^{-1}\left(F_{D}^{\alpha, q}(f) . \psi_{y}^{\alpha, q}\right)$ is in $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$ (resp. $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ ). Moreover, using the Plancherel formula and the relation (15), we obtain

$$
\begin{aligned}
\left\|T_{y}^{\alpha ; q} f\right\|_{2, \alpha, q} & =\left\|\left(F_{D}^{\alpha, q}\right)^{-1}\left(F_{D}^{\alpha, q}(f) \cdot \psi_{y}^{\alpha, q}\right)\right\|_{2, \alpha, q} \\
& =\left\|F_{D}^{\alpha, q}(f) \cdot \psi_{y}^{\alpha, q}\right\|_{2, \alpha, q} \leq \frac{4}{(q ; q)_{\infty}}\|f\|_{2, \alpha, q}
\end{aligned}
$$

4) Using the orthogonality relation (16), we obtain

$$
F_{D}^{\alpha, q}\left(\psi_{\lambda}^{\alpha, q}\right)(y)=\frac{2(1+q)^{\alpha} \Gamma_{q^{2}}(\alpha+1)}{(1-q)|\lambda y|^{\alpha+1}} \delta_{\lambda, y}, \quad \lambda, y \in \mathbb{R}_{q}
$$

witch implies that $T_{y}^{\alpha ; q}\left(\psi_{\lambda}^{\alpha, q}\right)(x)=\psi_{\lambda}^{\alpha, q}(x) \psi_{\lambda}^{\alpha, q}(y)$.
5) From the fact that for $f \in L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$ and $y \in \mathbb{R}_{q}, T_{y}^{\alpha ; q} f=\left(F_{D}^{\alpha, q}\right)^{-1}\left[F_{D}^{\alpha, q}(f) \cdot \psi_{y}^{\alpha, q}\right]$, we get

$$
F_{D}^{\alpha, q}\left(T_{y}^{\alpha ; q} f\right)(\lambda)=\left[F_{D}^{\alpha, q}(f) \cdot \psi_{y}^{\alpha, q}\right](\lambda)=\psi_{\lambda}^{\alpha, q}(y) \cdot F_{D}^{\alpha, q}(f)(\lambda)
$$

6) For $f \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, we have

$$
F_{D}^{\alpha, q}\left(\Lambda_{\alpha, q} T_{y}^{\alpha ; q} f\right)(\lambda)=i \lambda F_{D}^{\alpha, q}\left(T_{y}^{\alpha ; q} f\right)(\lambda)=i \lambda \psi_{\lambda}^{\alpha, q}(y) F_{D}^{\alpha, q}(f)(\lambda)
$$

and

$$
F_{D}^{\alpha, q}\left(T_{y}^{\alpha ; q} \Lambda_{\alpha, q} f\right)(\lambda)=\psi_{\lambda}^{\alpha, q}(y) F_{D}^{\alpha, q}\left(\Lambda_{\alpha, q} f\right)(\lambda)=i \lambda \psi_{\lambda}^{\alpha, q}(y) F_{D}^{\alpha, q}(f)(\lambda)
$$

The result follows, then, from the fact that $F_{D}^{\alpha, q}$ is an automorphism of $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$.

Definition 2. The $q$-convolution product is defined for $f, g \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ by:

$$
\begin{equation*}
f * g(x)=c_{\alpha, q} \int_{-\infty}^{\infty} T_{x}^{\alpha ; q} f(-y) g(y)|y|^{2 \alpha+1} d_{q} y \tag{27}
\end{equation*}
$$

In the following proposition, we present some of its properties.
Proposition 3. For $f, g \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, we have
i) $F_{D}^{\alpha, q}(f * g)=F_{D}^{\alpha, q}(f) \cdot F_{D}^{\alpha, q}(g)$.
ii) $f * g=g * f$.
iii) $(f * g) * h=f *(g * h)$.

Proof.
i) Let $f, g \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$. Then, with the help of the relation (15), we have for all $x, y \in \mathbb{R}_{q}$,

$$
\begin{aligned}
\left|T_{x}^{\alpha ; q} f(-y)\right| & \leq c_{\alpha, q} \int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}(f)(\lambda) \psi_{\lambda}^{\alpha, q}(x) \psi_{\lambda}^{\alpha, q}(-y)\right||\lambda|^{2 \alpha+1} d_{q} \lambda \\
& \leq c_{\alpha, q}\left(\frac{4}{(q ; q)_{\infty}}\right)^{2}\left\|F_{D}^{\alpha, q}(f)\right\|_{1, \alpha, q}
\end{aligned}
$$

So, since for $\lambda \in \mathbb{R}_{q}, \psi_{-\lambda}^{\alpha, q} \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, we obtain

$$
\begin{aligned}
& \int_{-\infty}^{\infty} \int_{-\infty}^{\infty}\left|T_{x}^{\alpha ; q} f(-y) g(y) \psi_{-\lambda}^{\alpha, q}(x) \| y\right|^{2 \alpha+1}|x|^{2 \alpha+1} d_{q} y d_{q} x \\
\leq & \frac{16 c_{\alpha, q}}{(q ; q)_{\infty}^{2}}\left\|F_{D}^{\alpha, q}(f)\right\|_{1, \alpha, q}\|g\|_{1, \alpha, q}\left\|\psi_{-\lambda}^{\alpha, q}\right\|_{1, \alpha, q} .
\end{aligned}
$$

Hence, using the Fubini's theorem and the properties of the generalized $q$ Dunkl translation, we get

$$
\begin{aligned}
F_{D}^{\alpha, q}(f * g)(\lambda) & =c_{\alpha, q} \int_{-\infty}^{\infty}(f * g)(x) \psi_{-\lambda}^{\alpha, q}(x)|x|^{2 \alpha+1} d_{q} x \\
& =c_{\alpha, q}^{2} \int_{-\infty}^{\infty}\left[\int_{-\infty}^{\infty} T_{x}^{\alpha ; q} f(-y) g(y)|y|^{2 \alpha+1} d_{q} y\right] \psi_{-\lambda}^{\alpha, q}(x)|x|^{2 \alpha+1} d_{q} x \\
& =c_{\alpha, q}^{2} \int_{-\infty}^{\infty}\left[\int_{-\infty}^{\infty} T_{-y}^{\alpha ; q} f(x) \psi_{-\lambda}^{\alpha, q}(x)|x|^{2 \alpha+1} d_{q} x\right] g(y)|y|^{2 \alpha+1} d_{q} y \\
& =c_{\alpha, q} \int_{-\infty}^{\infty} F_{D}^{\alpha, q}\left(T_{-y}^{\alpha ; q} f\right)(\lambda) g(y)|y|^{2 \alpha+1} d_{q} y \\
& =c_{\alpha, q} \int_{-\infty}^{\infty} F_{D}^{\alpha, q}(f)(\lambda) \psi_{\lambda}^{\alpha, q}(-y) g(y)|y|^{2 \alpha+1} d_{q} y \\
& =c_{\alpha, q} F_{D}^{\alpha, q}(f)(\lambda) \int_{-\infty}^{\infty} \psi_{-\lambda}^{\alpha, q}(y) g(y)|y|^{2 \alpha+1} d_{q} y \\
& =F_{D}^{\alpha, q}(f)(\lambda) F_{D}^{\alpha, q}(g)(\lambda) .
\end{aligned}
$$

ii) and iii) follows from i).

Proposition 4. Let $f$ and $g$ be in $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$. Then

1) $f * g \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$,
2) 

$$
\begin{equation*}
\int_{-\infty}^{\infty}|f * g(x)|^{2}|x|^{2 \alpha+1} d_{q} x=\int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}(f)(x)\right|^{2}\left|F_{D}^{\alpha, q}(g)(x)\right|^{2}|x|^{2 \alpha+1} d_{q} x \tag{28}
\end{equation*}
$$

Proof. The proof is a direct consequence of Theorem 2 and the fact that $F_{D}^{\alpha, q}(f * g)=F_{D}^{\alpha, q}(f) F_{D}^{\alpha, q}(g)$.

## 4. $q$-wavelet Transforms Associated with the $q$-Dunkl Operator

Definition 3. A q-wavelet, associated with the $q$-Dunkl operator, is a square $q$-integrable function $g$ on $\mathbb{R}_{q}$ satisfying the following admissibility condition

$$
\begin{equation*}
0<C_{\alpha, g}=\int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}(g)(a)\right|^{2} \frac{d_{q} a}{|a|}<\infty . \tag{29}
\end{equation*}
$$

## Remark 1.

1) For all $\lambda \in \mathbb{R}_{q}$, we have

$$
C_{\alpha, g}=\int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}(g)(a \lambda)\right|^{2} \frac{d_{q} a}{|a|} .
$$

2) Let $f$ be a nonzero function in $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$. Then $g=\Lambda_{\alpha, q} f$ is a $q$-wavelet, in $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ and we have

$$
C_{\alpha, g}=\int_{-\infty}^{\infty}|a|\left|F_{D}^{\alpha, q}(f)(a)\right|^{2} d_{q} a
$$

Proposition 5. Let $g \neq 0$ be a function in $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$ satisfying:
(1) $F_{D}^{\alpha, q}(g)$ is continuous at 0.

$$
\text { (2) } \exists \beta>0 \text { such that } F_{D}^{\alpha, q}(g)(x)-F_{D}^{\alpha, q}(g)(0)=O\left(x^{\beta}\right) \text {, as } x \rightarrow 0, x \in \mathbb{R}_{q} \text {. }
$$

Then, the admissibility condition (29) is equivalent to

$$
\begin{equation*}
F_{D}^{\alpha, q}(g)(0)=0 \tag{30}
\end{equation*}
$$

Proof. Assume that (29) is satisfied. If $F_{D}^{\alpha, q}(g)(0) \neq 0$, then there exist $p_{0} \in \mathbb{N}$ and $M>0$, such that

$$
\forall n \geq p_{0}, \quad\left|F_{D}^{\alpha, q}(g)\left( \pm q^{n}\right)\right| \geq M
$$

So, the $q$-integral in (29) would be equal to $\infty$.

- Conversely, assume that $F_{D}^{\alpha, q}(g)(0)=0$.

Since $g \neq 0$, we deduce from Theorem 2, that the first inequality in (29) holds.

On the other hand, from the assertion (2), there exist $n_{0} \in \mathbb{N}$ and $\epsilon>0$, such that for all $n \geq n_{0}$, we have

$$
\left|F_{D}^{\alpha, q}(g)\left( \pm q^{n}\right)\right| \leq \epsilon q^{n \beta} .
$$

Then using the definition of the $q$-integral and Theorem 2, we obtain

$$
\begin{aligned}
\int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}(g)(a)\right|^{2} \frac{d_{q} a}{|a|} & =(1-q) \sum_{n=-\infty}^{\infty}\left[\left|F_{D}^{\alpha, q}(g)\left(q^{n}\right)\right|^{2}+\left|F_{D}^{\alpha, q}(g)\left(-q^{n}\right)\right|^{2}\right] \\
& \leq \frac{\left\|F_{D}^{\alpha, q}(g)\right\|_{2, \alpha, q}^{2}}{q^{(2 \alpha+2) n_{0}}}+\frac{2(1-q)}{1-q^{2 \beta}} \epsilon^{2}
\end{aligned}
$$

This proves the second inequality of (29).

## Remark 2.

Using the relation (18), the continuity assumption in the previous proposition will certainly hold if $g$ is moreover in $L_{\alpha, q}^{1}\left(\mathbb{R}_{q}\right)$. Then (30) can be equivalently written as

$$
\int_{-\infty}^{\infty} g(x)|x|^{2 \alpha+1} d_{q} x=0
$$

which implies that $g$ must have sign changes on $\mathbb{R}_{q}$. It will also decay to 0 as $t$ tends to $\pm \infty$ (in $\mathbb{R}_{q}$ ). This explains the name " $q$-wavelet".

Proposition 6. For $a \in \mathbb{R}_{q}$ and $g \in L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$ (resp. $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ ), the function

$$
g_{a}: x \mapsto \frac{1}{|a|^{2 \alpha+2}} g\left(\frac{x}{a}\right)
$$

belongs to $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)\left(\right.$ resp. $\left.\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)\right)$ and we have

$$
\begin{equation*}
\left\|g_{a}\right\|_{2, \alpha, q}=\frac{1}{|a|^{\alpha+1}}\|g\|_{2, \alpha, q} \tag{31}
\end{equation*}
$$

and

$$
\begin{equation*}
F_{D}^{\alpha, q}\left(g_{a}\right)(\lambda)=F_{D}^{\alpha, q}(g)(a \lambda), \quad \lambda \in \widetilde{\mathbb{R}}_{q} \tag{32}
\end{equation*}
$$

Proof. The change of variables $u=\frac{x}{a}$ gives the result.

Proposition 7. Let $g$ be a q-wavelet, associated with he $q$-Dunkl operator, in $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$ (resp. $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ ). Then, for all $a \in \mathbb{R}_{q}$ and $b \in \widetilde{\mathbb{R}}_{q}$, the function $g_{(a, b), \alpha}$ defined by

$$
\begin{equation*}
g_{(a, b), \alpha}(x)=\sqrt{|a|} T_{b}^{\alpha ; q}\left(g_{a}\right)(x), \quad x \in \mathbb{R}_{q} \tag{33}
\end{equation*}
$$

is a q-wavelet associated with he $q$-Dunkl operator in $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)\left(\operatorname{resp} . \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)\right.$ ) and we have

$$
\begin{equation*}
C_{\alpha, g_{(a, b), \alpha}}=|a| \int_{-\infty}^{\infty}\left|\psi_{b}^{\alpha, q}\left(\frac{x}{a}\right)\right|^{2}\left|F_{D}^{\alpha, q}(g)(x)\right|^{2} \frac{d_{q} x}{|x|}, \tag{34}
\end{equation*}
$$

with $T_{b}^{\alpha ; q}$ is the generalized $q$-Dunkl translation operator defined by (24).
Proof. Using Proposition 2, Proposition 6 and the properties of the generalized $q$-Dunkl translation, we obtain $g_{(a, b), \alpha}$ is in $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)\left(\right.$ resp. $\left.\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)\right)$.
From Lemma 1 and the relations (26) and (32), we get for $a \in \mathbb{R}_{q}$ and $b \in \widetilde{\mathbb{R}}_{q}$,

$$
\begin{aligned}
C_{\alpha, g_{(a, b), \alpha}} & =|a| \int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}\left[T_{b}^{\alpha ; q}\left(g_{a}\right)\right](x)\right|^{2} \frac{d_{q} x}{|x|} \\
& =|a| \int_{-\infty}^{\infty}\left|\psi_{b}^{\alpha, q}\left(\frac{x}{a}\right)\right|^{2}\left|F_{D}^{\alpha, q}(g)(x)\right|^{2} \frac{d_{q} x}{|x|}
\end{aligned}
$$

Thus, since $g \neq 0$, we have, from the Plancherel theorem, $F_{D}^{\alpha, q}(g) \neq 0$ and

$$
0<C_{\alpha, g_{(a, b), \alpha}} \leq|a|\left(\frac{4}{(q ; q)_{\infty}}\right)^{2} \int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}(g)(x)\right|^{2} \frac{d_{q} x}{|x|}=\frac{16|a|}{(q ; q)_{\infty}^{2}} C_{\alpha, g}<+\infty
$$

Proposition 8. Let $g$ be a q-wavelet, associated with the $q$-Dunkl operator, in $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$. Then the mapping

$$
F:(a, b) \mapsto g_{(a, b), \alpha}
$$

is continuous from $\mathbb{R}_{q} \times \widetilde{\mathbb{R}}_{q}$ into $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$, via the induced topology on $\mathbb{R}_{q} \times \widetilde{\mathbb{R}}_{q}$ by that of $\mathbb{R} \times \mathbb{R}$.

Proof. It is clear, from the previous proposition, that $F$ is a mapping from $\mathbb{R}_{q} \times \widetilde{\mathbb{R}}_{q}$ into $L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$ and it is continuous on $\mathbb{R}_{q} \times \mathbb{R}_{q}$, since every element of
$\mathbb{R}_{q} \times \mathbb{R}_{q}$ is an isolated point.
Now, let $a \in \mathbb{R}_{q}$. For $b \in \widetilde{\mathbb{R}}_{q}$, we have

$$
\begin{aligned}
\|F(a, b)-F(a, 0)\|_{2, \alpha, q}^{2} & =|a|\left\|T_{b}^{\alpha ; q}\left(g_{a}\right)-g_{a}\right\|_{2, \alpha, q}^{2} \\
& =|a|\left\|F_{D}^{\alpha, q}\left(T_{b}^{\alpha ; q}\left(g_{a}\right)-g_{a}\right)\right\|_{2, \alpha, q}^{2} \\
& =|a| \int_{-\infty}^{\infty}\left|1-\psi_{b}^{\alpha, q}(x)\right|^{2}\left|F_{D}^{\alpha, q}\left(g_{a}\right)\right|^{2}(x)|x|^{2 \alpha+1} d_{q} x .
\end{aligned}
$$

Using the relation (15), the fact that $F_{D}^{\alpha, q}\left(g_{a}\right) \in L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$ and the Lebesgue theorem we obtain

$$
\lim _{\substack{b \rightarrow 0 \\ b \in \tilde{\mathbb{R}}_{q}}}\|F(a, b)-F(a, 0)\|_{2, \alpha, q}=0
$$

Definition 4. Let $g$ be a q-wavelet, associated with the $q$-Dunkl operator, in $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$. We define the continuous $q$-wavelet transform associated with the $q$ Dunkl operator for $f \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, by

$$
\begin{equation*}
\Psi_{q, g}^{\alpha}(f)(a, b)=c_{\alpha, q} \int_{-\infty}^{\infty} f(x) \overline{g_{(a, b), \alpha}}(-x)|x|^{2 \alpha+1} d_{q} x, \quad a \in \mathbb{R}_{q}, \quad b \in \widetilde{\mathbb{R}}_{q} \tag{35}
\end{equation*}
$$

Remark that (35) is equivalent to

$$
\begin{aligned}
\Psi_{q, g}^{\alpha}(f)(a, b) & =\sqrt{|a|} f * \overline{g_{a}}(b) \\
& =\sqrt{|a|} F_{D}^{\alpha, q}\left[F_{D}^{\alpha, q}\left(f * \overline{g_{a}}\right)\right](-b) \\
& =\sqrt{|a|} F_{D}^{\alpha, q}\left[F_{D}^{\alpha, q}(f) \cdot F_{D}^{\alpha, q}\left(\overline{g_{a}}\right)\right](-b) \\
& =\sqrt{|a|} c_{\alpha, q} \int_{-\infty}^{\infty} F_{D}^{\alpha, q}(f)(x) \cdot F_{D}^{\alpha, q}(\bar{g})(a x) \psi_{b}^{\alpha, q}(x)|x|^{2 \alpha+1} d_{q} x,
\end{aligned}
$$

where $c_{\alpha, q}$ is given by (10).
The following propositions give some properties of $\Psi_{q, g}^{\alpha}$.
Proposition 9. Let $g$ be a q-wavelet, associated with the $q$-Dunkl operator, in $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ and $f \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$. Then
i) For all $a \in \mathbb{R}_{q}$ and $b \in \widetilde{\mathbb{R}}_{q}$, we have

$$
\begin{equation*}
\left|\Psi_{q, g}^{\alpha}(f)(a, b)\right| \leq \frac{4 c_{\alpha, q}}{|a|^{\alpha+\frac{1}{2}}(q ; q)_{\infty}}\|f\|_{2, \alpha, q}\|g\|_{2, \alpha, q} \tag{36}
\end{equation*}
$$

ii) For all $a \in \mathbb{R}_{q}$, the mapping $b \mapsto \Psi_{q, g}^{\alpha}(f)(a, b)$ is continuous on $\widetilde{\mathbb{R}}_{q}$, via the induced topology on $\widetilde{\mathbb{R}}_{q}$ by that of $\mathbb{R}$, and we have

$$
\begin{equation*}
\lim _{b \rightarrow \infty} \Psi_{q, g}^{\alpha}(f)(a, b)=0 \tag{37}
\end{equation*}
$$

Proof. i) Using the properties of the generalized $q$-Dunkl translation operator, the Cauchy-Schwartz inequality and Lemma 1, we obtain for $a \in \mathbb{R}_{q}$ and $b \in \widetilde{\mathbb{R}}_{q}$,

$$
\begin{aligned}
\left|\Psi_{q, g}^{\alpha}(f)(a, b)\right| & =\left.c_{\alpha, q}\left|\int_{-\infty}^{\infty} f(x) \overline{g_{(a, b), \alpha}}(-x)\right| x\right|^{2 \alpha+1} d_{q} x \mid \\
& \leq \sqrt{|a|} c_{\alpha, q} \int_{-\infty}^{\infty}\left|f(x) \| T_{b}^{\alpha ; q} g_{a}(-x)\right||x|^{2 \alpha+1} d_{q} x \\
& \leq \frac{4 c_{\alpha, q}}{|a|^{\alpha+\frac{1}{2}}(q ; q)_{\infty}}\|f\|_{2, \alpha, q}\|g\|_{2, \alpha, q} .
\end{aligned}
$$

ii) Since every element of $\mathbb{R}_{q}$ is an isolated point, it is sufficient to prove the continuity at 0 . For $b \in \widetilde{\mathbb{R}}_{q}$, we have

$$
\Psi_{q, g}^{\alpha}(f)(a, b)=\sqrt{|a|} F_{D}^{\alpha, q}\left[F_{D}^{\alpha, q}(f) \cdot F_{D}^{\alpha, q}\left(\overline{g_{a}}\right)\right](-b)
$$

Since $f, g \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, then from Theorem 2, we have $F_{D}^{\alpha, q}(f)$ and $F_{D}^{\alpha, q}\left(\overline{g_{a}}\right)$ are in $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ and the product $F_{D}^{\alpha, q}(f) \cdot F_{D}^{\alpha, q}\left(\overline{g_{a}}\right)$ is in $L_{\alpha, q}^{1}\left(\mathbb{R}_{q}\right)$. Thus, using the relation (15), the Lebesgue theorem, gives

$$
\begin{aligned}
\lim _{\substack{b \rightarrow 0 \\
b \in \mathbb{R}_{q}}} \Psi_{q, g}^{\alpha}(f)(a, b) & =\lim _{\substack{b \rightarrow 0 \\
b \in \mathbb{R}_{q}}} \sqrt{|a|} c_{\alpha, q} \int_{-\infty}^{\infty} F_{D}^{\alpha, q}(f)(x) \cdot F_{D}^{\alpha, q}\left(\overline{g_{a}}\right)(x) \psi_{b}^{\alpha, q}(x) d_{q} x \\
& =\Psi_{q, g}^{\alpha}(f)(a, 0)
\end{aligned}
$$

Which proves the continuity of $\Psi_{q, g}^{\alpha}(f)(a,$.$) at 0$.
Finally, the relation (18) implies that

$$
\Psi_{q, g}^{\alpha}(a, b)=\sqrt{|a|} F_{D}^{\alpha, q}\left[F_{D}^{\alpha, q}(f) \cdot F_{D}^{\alpha, q}\left(\overline{g_{a}}\right)\right](-b)
$$

tends to 0 when $b$ tends to $\infty$.
Let us now establish a Plancherel and a Parseval formulas for $\Psi_{q, g}^{\alpha}$.

Theorem 3. Let $g \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ be a $q$-wavelet, associated with the $q$-Dunkl operator.
i) Plancherel formula

For $f \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, we have

$$
\begin{equation*}
\frac{1}{C_{\alpha, g}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty}\left|\Psi_{q, g}^{\alpha}(f)(a, b)\right|^{2}|b|^{2 \alpha+1} \frac{d_{q} a d_{q} b}{|a|^{2}}=\|f\|_{2, \alpha, q}^{2} \tag{38}
\end{equation*}
$$

## ii) Parseval formula

For $f_{1}, f_{2} \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, we have

$$
\begin{equation*}
\int_{-\infty}^{\infty} f_{1}(x) \bar{f}_{2}(x)|x|^{2 \alpha+1} d_{q} x=\frac{1}{C_{\alpha, g}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Psi_{q, g}^{\alpha}\left(f_{1}\right)(a, b) \overline{\Psi_{q, g}^{\alpha}\left(f_{2}\right)}(a, b)|b|^{2 \alpha+1} \frac{d_{q} a d_{q} b}{|a|^{2}} \tag{39}
\end{equation*}
$$

Proof. The use of the Fubini's theorem, Theorem 2, the relations (28) and (32) gives

$$
\begin{aligned}
& \int_{-\infty}^{\infty} \int_{-\infty}^{\infty}\left|\Psi_{q, g}^{\alpha}(f)(a, b)\right|^{2}|b|^{2 \alpha+1} \frac{d_{q} a d_{q} b}{|a|^{2}} \\
= & \int_{-\infty}^{\infty}\left(\int_{-\infty}^{\infty}\left|f * \overline{g_{a}}(b)\right|^{2}|b|^{2 \alpha+1} d_{q} b\right) \frac{d_{q} a}{|a|} \\
= & \int_{-\infty}^{\infty}\left(\int_{-\infty}^{\infty}\left|\left[F_{D}^{\alpha, q}(f) F_{D}^{\alpha, q}\left(\overline{g_{a}}\right)\right](b)\right|^{2}|b|^{2 \alpha+1} d_{q} b\right) \frac{d_{q} a}{|a|} \\
= & \int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}(f)(b)\right|^{2}|b|^{2 \alpha+1}\left(\int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}(g)(a b)\right|^{2} \frac{d_{q} a}{|a|}\right) d_{q} b \\
= & C_{\alpha, g} \int_{-\infty}^{\infty}\left|F_{D}^{\alpha, q}(f)(b)\right|^{2}|b|^{2 \alpha+1} d_{q} b=C_{\alpha, g}\|f\|_{2, \alpha, q}^{2} .
\end{aligned}
$$

ii) The result follows from (38).

Theorem 4. Let $g \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ be a $q$-wavelet, associated with the $q$-Dunkl operator. Then for $f \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, we have

$$
\begin{equation*}
f(x)=\frac{c_{\alpha, q}}{C_{\alpha, g}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Psi_{q, g}^{\alpha}(f)(a, b) g_{(a, b), \alpha}(-x)|b|^{2 \alpha+1} \frac{d_{q} a d_{q} b}{|a|^{2}}, \quad x \in \mathbb{R}_{q} \tag{40}
\end{equation*}
$$

Proof. Let $x \in \mathbb{R}_{q}$ and put $h=\delta_{x}$. It is easy to see that $h \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$. According to the relation (39) of the previous theorem and the definition of $\Psi_{q, g}^{\alpha}$ and the $q$-Jackson integral, we have,

$$
\begin{aligned}
& (1-q)|x|^{2 \alpha+2} f(x)=\int_{-\infty}^{\infty} f(t) \bar{h}(t)|t|^{2 \alpha+1} d_{q} t \\
= & \frac{1}{C_{\alpha, g}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Psi_{q, g}^{\alpha}(f)(a, b) \overline{\Psi_{q, g}^{\alpha}}(h)(a, b)|b|^{2 \alpha+1} \frac{d_{q} a d_{q} b}{|a|^{2}} \\
= & \frac{c_{\alpha, q}}{C_{\alpha, g}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Psi_{q, g}^{\alpha}(f)(a, b)\left(\int_{-\infty}^{\infty} \bar{h}(t) g_{(a, b), \alpha}(-t)|t|^{2 \alpha+1} d_{q} t\right)|b|^{2 \alpha+1} \frac{d_{q} a d_{q} b}{|a|^{2}} \\
= & (1-q)|x|^{2 \alpha+2} \frac{c_{\alpha, q}}{C_{\alpha, g}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Psi_{q, g}^{\alpha}(f)(a, b) g_{(a, b), \alpha}(-x)|b|^{2 \alpha+1} \frac{d_{q} a d_{q} b}{|a|^{2}}
\end{aligned}
$$

## 5. Inversion Formulas for the $q$-Dunkl Intertwining Operator and its Dual

In the what follows, we will need the following spaces:

- $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)=\left\{f \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right): \int_{-\infty}^{+\infty} f(x) x^{k}|x|^{2 \alpha+1} d_{q} x=0, k=0,1, \ldots\right\}$.
- $\mathcal{S}_{q}^{0}\left(\mathbb{R}_{q}\right)=\left\{f \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right): \quad \partial_{q}^{k} f(0)=0, \quad k=0,1, \ldots\right\}$.

We recall that the $q$-Dunkl intertwining operator $V_{\alpha, q}$ is defined on $\mathcal{E}_{q}\left(\mathbb{R}_{q}\right)$ by (see [1])

$$
\begin{equation*}
V_{\alpha, q}(f)(x)=\frac{(1+q)}{2} \frac{\Gamma_{q^{2}}(\alpha+1)}{\Gamma_{q^{2}}\left(\frac{1}{2}\right) \Gamma_{q^{2}}\left(\alpha+\frac{1}{2}\right)} \int_{-1}^{1} \frac{\left(t^{2} q^{2} ; q^{2}\right)_{\infty}}{\left(t^{2} q^{2 \alpha+1} ; q^{2}\right)_{\infty}}(1+t) f(x t) d_{q} t \tag{41}
\end{equation*}
$$

The dual operator of $V_{\alpha, q}$ is defined on $\mathcal{D}_{q}\left(\mathbb{R}_{q}\right)$ by (see [1])

$$
\begin{equation*}
\left({ }^{t} V_{\alpha, q}\right)(f)(t)=\frac{(1+q)^{-\alpha+\frac{1}{2}}}{2 \Gamma_{q^{2}}\left(\alpha+\frac{1}{2}\right)} \int_{|x| \geq q|t|} \frac{\left(\left(\frac{t}{x}\right)^{2} q^{2} ; q^{2}\right)_{\infty}}{\left(\left(\frac{t}{x}\right)^{2} q^{2 \alpha+1} ; q^{2}\right)_{\infty}}\left(1+\frac{t}{x}\right) f(x) \frac{|x|^{2 \alpha+1}}{x} d_{q} x \tag{42}
\end{equation*}
$$

These two operators satisfy the following properties:
Proposition 10. i) $V_{\alpha, q}\left(e\left(-i \lambda x ; q^{2}\right)\right)=\psi_{-\lambda}^{\alpha, q}(x), \lambda, x \in \mathbb{R}_{q}$.
ii) For $f \in \mathcal{E}_{q}\left(\mathbb{R}_{q}\right)$ and $g \in \mathcal{D}_{q}\left(\mathbb{R}_{q}\right)$

$$
\begin{equation*}
c_{\alpha, q} \int_{-\infty}^{+\infty} V_{\alpha, q}(f)(x) g(x)|x|^{2 \alpha+1} d_{q} x=\frac{(1+q)^{\frac{1}{2}}}{2 \Gamma_{q^{2}}\left(\frac{1}{2}\right)} \int_{-\infty}^{+\infty} f(t)\left({ }^{t} V_{\alpha, q}\right)(g)(t) d_{q} t . \tag{43}
\end{equation*}
$$

iii) $V_{\alpha, q}$ and ${ }^{t} V_{\alpha, q}$ verify the following transmutation relations

$$
\begin{gather*}
\Lambda_{\alpha, q} V_{\alpha, q}(f)=V_{\alpha, q}\left(\partial_{q} f\right), \quad V_{\alpha, q}(f)(0)=f(0), \quad f \in \mathcal{E}_{q}\left(\mathbb{R}_{q}\right),  \tag{44}\\
\partial_{q}\left({ }^{t} V_{\alpha, q}\right)(f)=\left({ }^{t} V_{\alpha, q}\right)\left(\Lambda_{\alpha, q}\right)(f), \quad f \in \mathcal{D}_{q}\left(\mathbb{R}_{q}\right) . \tag{45}
\end{gather*}
$$

iv) The $q$-Dunkl transform and the $q^{2}$-analogue Fourier transform $\mathcal{F}_{q}$, studied in ([11], [12]), are linked by the following relation (see [1]):

$$
\begin{equation*}
\forall f \in \mathcal{D}_{q}\left(\mathbb{R}_{q}\right), \quad F_{D}^{\alpha, q}(f)=\mathcal{F}_{q} \circ \quad{ }^{t} V_{\alpha, q}(f) \tag{46}
\end{equation*}
$$

where

$$
\begin{equation*}
\mathcal{F}_{q}(f)(x)=F_{D}^{-\frac{1}{2}, q}(f)(x)=\frac{(1+q)^{\frac{1}{2}}}{2 \Gamma_{q^{2}}\left(\frac{1}{2}\right)} \int_{-\infty}^{\infty} f(t) e\left(-i t x ; q^{2}\right) d_{q} t \tag{47}
\end{equation*}
$$

We state the following results, useful in the sequel.
Theorem 5. The $q^{2}$-analogue Fourier transform $\mathcal{F}_{q}$ is an isomorphism from $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$ into $\mathcal{S}_{q}^{0}\left(\mathbb{R}_{q}\right)$.
Proof. The result follows from the fact that $\partial_{q} e\left(-i x ; q^{2}\right)=-i e\left(-i x ; q^{2}\right)$.
Similarly, we have the following result.
Theorem 6. The $q$-Dunkl transform $F_{D}^{\alpha, q}$ is an isomorphism from $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$ into $\mathcal{S}_{q}^{0}\left(\mathbb{R}_{q}\right)$.

Proof. On the one hand, we have for all $k \in \mathbb{N}$,

$$
\partial_{q}^{k} F_{D}^{\alpha, q}(f)(\lambda)=\int_{-\infty}^{+\infty} f(x)|x|^{2 \alpha+1} \partial_{q, \lambda}^{k}\left[\psi_{-\lambda}^{\alpha, q}(x)\right] d_{q} x
$$

On the other hand, from the relation (14), we have
$\partial_{q, \lambda}^{k}\left[\psi_{-\lambda}^{\alpha, q}\right](x)=\frac{(1+q) \Gamma_{q^{2}}(\alpha+1)}{2 \Gamma_{q^{2}}\left(\frac{1}{2}\right) \Gamma_{q^{2}}\left(\alpha+\frac{1}{2}\right)}(-i)^{k} x^{k} \int_{-1}^{1} \frac{\left(t^{2} q^{2} ; q^{2}\right)_{\infty}}{\left(t^{2} q^{2 \alpha+1} ; q^{2}\right)_{\infty}}(1+t) t^{k} e\left(-i \lambda x t, q^{2}\right) d_{q} t$, which gives the result.

Corollary 1. The operator ${ }^{t} V_{\alpha, q}$ is an isomorphism from $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$ into $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$.

Proof. We deduce the result from the relation $F_{D}^{\alpha, q}=\mathcal{F}_{q} \circ{ }^{t} V_{\alpha, q}$ and Theorems 5 and 6.

Proposition 11. For $f$ in $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$ (resp. $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$ ) and $g$ in $\mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$ the function $f *_{q} g$ (resp. $f * g$ ) belongs to $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$ (resp. $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$ ).
where "*${ }_{q}$ " is the $q$-convolution product associated with the operator $\partial_{q}$ studied in [11].

Proof. The result follows from Theorem 5 (resp. 6) and the fact that $f *_{q} g=\mathcal{F}_{q}^{-1}\left(\mathcal{F}_{q}(f) . \mathcal{F}_{q}(g)\right.$ ) (resp. $f * g=\left(F_{D}^{\alpha, q}\right)^{-1}\left(F_{D}^{\alpha, q}(f) . F_{D}^{\alpha, q}(g)\right)$.

Using the Taylor formula for the Jackson's $q$-derivative (see [3, 2]), we provide in the following lemma a Taylor formula for the operator $\partial_{q}$.

Lemma 2. Let $f$ be a function $N$ times continuously $q$-differentiable on $\widetilde{\mathbb{R}}_{q}$, $N \in \mathbb{N}$. Then,

$$
f(x)=\sum_{n=0}^{N} q^{\left(E\left(\frac{n+1}{2}\right)\right)^{2}} \frac{\partial_{q}^{n} f(0)}{[n]_{q}!} x^{n}+\frac{x^{N}}{[N]_{q}!} \int_{0}^{1}(t q ; q)_{N} \quad H_{q, N+1}(f)(x t) d_{q} t
$$

where for $n \in \mathbb{N}, E\left(\frac{n+1}{2}\right)$ is the integer part of $\frac{n+1}{2}$ and $H_{q, n}$ is the operator defined by

$$
H_{q, n}(f)(t)=q^{a_{n}} \partial_{q}^{n} f_{o}\left(t q^{E\left(1+\frac{n}{2}\right)}\right)+q^{b_{n}} \partial_{q}^{n} f_{e}\left(t q^{E\left(\frac{n+1}{2}\right)}\right)
$$

with $f_{o}$ and $f_{e}$ are respectively the odd and the even parts of $f$,
$a_{n}=\left\{\begin{array}{ll}\frac{n(n+2)}{4}, & \text { if } n \text { is even, } \\ \frac{(n+1)^{2}}{4}, & \text { if } n \text { is odd, }\end{array} \quad\right.$ and $\quad b_{n}=\left\{\begin{array}{lll}\frac{n^{2}}{4}, & \text { if } n \text { is even, } \\ \frac{(n-1)(n+1)}{4}, & \text { if } n \text { is odd. } .\end{array}\right.$
Proposition 12. The operator $K_{\alpha, q, 1}$ defined by

$$
K_{\alpha, q, 1}(f)=\frac{\Gamma_{q^{2}}(1 / 2)}{(1+q)^{(\alpha+1 / 2)} \Gamma_{q^{2}}(\alpha+1)} \mathcal{F}_{q}^{-1}\left(|\lambda|^{2 \alpha+1} \mathcal{F}_{q}(f)\right)
$$

is an isomorphism from $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$ onto itself.

Proof. Using the previous lemma, one can prove that the multiplication operator $f \mapsto \frac{\Gamma_{q^{2}}(1 / 2)}{(1+q)^{(\alpha+1 / 2)} \Gamma_{q^{2}}(\alpha+1)}|\lambda|^{2 \alpha+1} f \quad$ is an isomorphism from $\mathcal{S}_{q}^{0}\left(\mathbb{R}_{q}\right)$ onto itself, its inverse is given by $f \mapsto \frac{(1+q)^{(\alpha+1 / 2)} \Gamma_{q^{2}}(\alpha+1)}{\Gamma_{q^{2}}(1 / 2)|\lambda|^{2 \alpha+1}} f$. The result follows, then, from Theorem 5.

Proposition 13. The operator $K_{\alpha, q, 2}$ defined by

$$
K_{\alpha, q, 2}(f)(x)=\frac{\Gamma_{q^{2}}(1 / 2)}{(1+q)^{(\alpha+1 / 2)} \Gamma_{q^{2}}(\alpha+1)}\left(F_{D}^{\alpha, q}\right)^{-1}\left(|\lambda|^{2 \alpha+1} F_{D}^{\alpha, q}(f)\right)(x)
$$

is an isomorphism from $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$ onto itself.
Proof. From the relation $F_{D}^{\alpha, q}=\mathcal{F}_{q} \circ{ }^{t} V_{\alpha, q}$ and the definition of $K_{\alpha, q, 1}$, we have for all $f \in \mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$,

$$
\begin{equation*}
K_{\alpha, q, 2}=\left({ }^{t} V_{\alpha, q}\right)^{-1} \circ K_{\alpha, q, 1} \circ{ }^{t} V_{\alpha, q} . \tag{48}
\end{equation*}
$$

We deduce the result from Proposition 12 and Corollary 1.

## Proposition 14.

i) For all $f \in \mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$ and $g \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, we have

$$
K_{\alpha, q, 1}\left(f *_{q} g\right)=K_{\alpha, q, 1}(f) *_{q} g
$$

ii) For all $f \in \mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$ and $g \in \mathcal{S}_{q}\left(\mathbb{R}_{q}\right)$, we have

$$
K_{\alpha, q, 2}(f * g)=K_{\alpha, q, 2}(f) * g
$$

Proof. The result follows from the properties of the $q$-convolution product and the definitions of $K_{\alpha, q, 1}$ and $K_{\alpha, q, 2}$.

Theorem 7. For all $f \in \mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$, we have the following inversion formulas for the operator $V_{\alpha, q}$

$$
\begin{equation*}
f=V_{\alpha, q} \circ K_{\alpha, q, 1} \circ{ }^{t} V_{\alpha, q}(f) \tag{49}
\end{equation*}
$$

and

$$
\begin{equation*}
f=V_{\alpha, q} \circ t V_{\alpha, q} \circ K_{\alpha, q, 2}(f) \tag{50}
\end{equation*}
$$

Proof. Using the properties of the operator $V_{\alpha, q}$, studied in [1], Theorem 1 and relation (46), we obtain for $x \in \widetilde{\mathbb{R}}_{q}$,

$$
\begin{aligned}
f(x) & =c_{\alpha, q} \int_{-\infty}^{+\infty} F_{D}^{\alpha, q}(f)(\lambda) \psi_{\lambda}^{\alpha, q}(x) \cdot|\lambda|^{2 \alpha+1} d_{q} \lambda \\
& =V_{\alpha, q}\left[c_{\alpha, q} \int_{-\infty}^{\infty} F_{D}^{\alpha, q}(f)(\lambda) e\left(i \lambda_{\bullet} ; q^{2}\right)|\lambda|^{2 \alpha+1} d_{q} \lambda\right](x) \\
& =V_{\alpha, q}\left\{\frac{c_{\alpha, q}}{c_{-1 / 2, q}} \mathcal{F}_{q}^{-1}\left[|\lambda|^{2 \alpha+1} \mathcal{F}_{q} \circ{ }^{t} V_{\alpha, q}(f)\right]\right\}(x) \\
& =V_{\alpha, q} \circ K_{\alpha, q, 1} \circ{ }^{t} V_{\alpha, q}(f)(x) .
\end{aligned}
$$

We deduce the second from the first relation and the the relation (48).
Corollary 2. The operator $V_{\alpha, q}$ is an isomorphism from $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$ into $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$.

Proof. We deduce the result from Proposition 12, Corollary 1 and the relation (49).

Similarly, we have the following result.
Theorem 8. For all $f \in \mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$, we have the following inversion formulas for the operator ${ }^{t} V_{\alpha, q}$

$$
\begin{equation*}
f={ }^{t} V_{\alpha, q} \circ V_{\alpha, q} \circ K_{\alpha, q, 1}(f) \tag{51}
\end{equation*}
$$

and

$$
\begin{equation*}
f={ }^{t} V_{\alpha, q} \circ K_{\alpha, q, 2} \circ V_{\alpha, q}(f) \tag{52}
\end{equation*}
$$

Proof. For $f \in \mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$, Corollary 1 (resp. 2) implies that ${ }^{t} V_{\alpha, q}^{-1}(f)$ (resp. $V_{\alpha, q}(f)$ ) belongs to $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$. Then by writing the relation (49) (resp. (50)) for ${ }^{t} V_{\alpha, q}^{-1}(f)$ (resp. $V_{\alpha, q}(f)$ ), we obtain the result.

Corollary 3. i) For all $f, g \in \mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$, we have

$$
\begin{equation*}
{ }^{t} V_{\alpha, q}(f * g)={ }^{t} V_{\alpha, q}(f) *{ }_{q}{ }^{t} V_{\alpha, q}(g) \tag{53}
\end{equation*}
$$

ii) For all $f, g \in \mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$ we have

$$
\begin{equation*}
V_{\alpha, q}\left(f *_{q} g\right)=V_{\alpha, q}(f) *{ }^{t} V_{\alpha, q}^{-1}(g) \tag{54}
\end{equation*}
$$

## 6. Inversion of the $q$-Dunkl Intertwining Operator and of its Dual using Wavelets

In this section, we assume that the reader is familiar with the notions and notations presented in [4], where the authors studied the particular case $\alpha=$ $-\frac{1}{2}$. In particular, we recall the following notations

$$
\begin{gathered}
H_{a}(f)(x)=\frac{1}{\sqrt{|a|}} f\left(\frac{x}{a}\right), \quad C_{g}=\int_{-\infty}^{\infty}\left|\mathcal{F}_{q}(g)\right|^{2}(a) \frac{d_{q} a}{|a|} \\
g_{a, b}=g_{(a, b),-1 / 2} \quad \text { and } \quad \Phi_{q, g}=\Psi_{q, g}^{-1 / 2}
\end{gathered}
$$

We begin by the following useful and easily verified result.
Proposition 15. For all $a \in \mathbb{R}_{q}$ and all $g \in L_{\alpha, q}^{2}\left(\mathbb{R}_{q}\right)$, we have

$$
\begin{aligned}
g_{a}=\frac{1}{|a|^{2 \alpha+3 / 2}} H_{a}(g) & =\frac{1}{\sqrt{|a|}}\left(F_{D}^{\alpha, q}\right)^{-1} \circ H_{a^{-1}} \circ F_{D}^{\alpha, q}(g) \\
& =\frac{1}{\sqrt{|a|}}{ }^{t} V_{\alpha, q}^{-1} \circ H_{a} \circ{ }^{t} V_{\alpha, q}(g) .
\end{aligned}
$$

Proposition 16. Let $g \in \mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$ be a $q$-wavelet, associated with the $q$-Dunkl operator. Then, for all $f$ in $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$, we have

$$
\begin{equation*}
\Psi_{q, g}^{\alpha}(f)(a, .)={ }^{t} V_{\alpha, q}^{-1}\left[\Phi_{q, \quad{ }^{t} V_{\alpha, q}(g)}\left({ }^{t} V_{\alpha, q}(f)\right)(a, .)\right], \quad a \in \mathbb{R}_{q} . \tag{55}
\end{equation*}
$$

Proof. Let $a \in \mathbb{R}_{q}$, from the properties of the continuous $q$-wavelet transform (see [4]), the relation (53) and Proposition (15), we have

$$
\begin{aligned}
\Psi_{q, g}^{\alpha}(f)(a, .) & =\sqrt{|a|} f * \overline{g_{a}}=\sqrt{|a|}{ }^{t} V_{\alpha, q}^{-1}\left[{ }^{t} V_{\alpha, q}(f) *_{q}{ }^{t} V_{\alpha, q}\left(\overline{g_{a}}\right)\right] \\
& ={ }^{t} V_{\alpha, q}^{-1}\left[{ }^{t} V_{\alpha, q}(f) *_{q} \overline{H_{a} \circ{ }^{t} V_{\alpha, q}(g)}\right] \\
& ={ }^{t} V_{\alpha, q}^{-1}\left[\Phi_{q,}{ }^{t} V_{\alpha, q}(g)\left({ }^{t} V_{\alpha, q}(f)\right)(a, .)\right] .
\end{aligned}
$$

Theorem 9. Let $g \in \mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$ be a q-wavelet, associated with the $q$-Dunkl operator. Then,

1) For all $f$ in $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$, we have for $a \in \mathbb{R}_{q}$ and $b \in \widetilde{\mathbb{R}}_{q}$,

$$
\begin{equation*}
\Psi_{q, g}^{\alpha}(f)(a, b)=V_{\alpha, q}\left[\Phi_{q,} \quad{ }^{t} V_{\alpha, q}(g)\left(V_{\alpha, q}^{-1}(f)\right)(a, .)\right](b), \tag{56}
\end{equation*}
$$

2) For all $f$ in $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$, we have for $a \in \mathbb{R}_{q}$ and $b \in \widetilde{\mathbb{R}}_{q}$,

$$
\begin{equation*}
\Phi_{q,}{ }^{t} V_{\alpha, q}(g)(f)(a, b)={ }^{t} V_{\alpha, q}\left[\Psi_{q, g}^{\alpha}\left({ }^{t} V_{\alpha, q}^{-1}(f)\right)(a, .)\right](b) . \tag{57}
\end{equation*}
$$

Proof. We deduce the result from Proposition 15, Corollary 3, the properties of the continuous $q$-wavelet transform (see [4]) and the relation (55).
Proposition 17. 1) If $g$ is a $q$-wavelet (associated with the operator $\partial_{q}$ ) in $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$, then $K_{\alpha, q, 1}(g)$ is a q-wavelet in $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$ and we have

$$
\begin{equation*}
K_{\alpha, q, 1} \circ H_{a}(g)=\frac{1}{|a|^{2 \alpha+1}} H_{a} \circ K_{\alpha, q, 1}(g), \quad a \in \mathbb{R}_{q} \tag{58}
\end{equation*}
$$

2) If $g$ is a $q$-wavelet, associated with the $q$-Dunkl operator, in $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$, then $K_{\alpha, q, 2}(g)$ is a $q$-wavelet, associated with the $q$-Dunkl operator, in $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$ and we have

$$
\begin{equation*}
K_{\alpha, q, 2}\left(g_{a}\right)=\frac{1}{|a|^{2 \alpha+1}}\left(K_{\alpha, q, 2}(g)\right)_{a}, \quad a \in \mathbb{R}_{q} \tag{59}
\end{equation*}
$$

Proof. 1) Let $g$ be a $q$-wavelet in $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$. From the definition of $K_{\alpha, q, 1}$, we have for $\lambda \in \mathbb{R}_{q}, \mathcal{F}_{q}\left(K_{\alpha, q, 1}(g)\right)(\lambda)=\frac{\Gamma_{q^{2}}(1 / 2)}{(1+q)^{(\alpha+1 / 2)} \Gamma_{q^{2}}(\alpha+1)}|\lambda|^{2 \alpha+1} \mathcal{F}_{q}(g)(\lambda)$.
Proposition 5 of [4], implies that $K_{\alpha, q, 1}(g)$ is a $q$-wavelet. On the other hand, using the fact $\mathcal{F}_{q} \circ H_{a}=H_{a^{-1}} \circ \mathcal{F}_{q}, \quad a \in \mathbb{R}_{q}$ and the above equality, we obtain
$\mathcal{F}_{q}\left(H_{a} \circ K_{\alpha, q, 1}(g)\right)(\lambda)=|a|^{2 \alpha+1} \frac{\Gamma_{q^{2}}(1 / 2)}{(1+q)^{(\alpha+1 / 2)} \Gamma_{q^{2}}(\alpha+1)}|\lambda|^{2 \alpha+1} \mathcal{F}_{q}\left(H_{a}(g)\right)(\lambda)$,
which gives the result.
2) The same way of 1 ) leads to the result.

Theorem 10. Let $g$ be a q-wavelet, associated with the $q$-Dunkl operator, in $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$. Then for $a \in \mathbb{R}_{q}$ and $b \in \mathbb{R}_{q}$, we have:

1) For all $f$ in $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$,

$$
\begin{equation*}
\Psi_{q, g}^{\alpha}(f)(a, b)=\frac{1}{|a|^{2 \alpha+1}} V_{\alpha, q}\left[\Phi_{q, K_{\alpha, q, 1} \circ} \quad{ }^{t} V_{\alpha, q}(g)\left({ }^{t} V_{\alpha, q}(f)\right)(a, .)\right](b) ; \tag{60}
\end{equation*}
$$

2) For all $f$ in $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$,

$$
\begin{equation*}
\Phi_{q,}{ }^{t} V_{\alpha, q}(g)(f)(a, b)=\frac{1}{|a|^{2 \alpha+1}}{ }^{t} V_{\alpha, q}\left[\Psi_{q, K_{\alpha, q, 2}(g)}^{\alpha}\left(V_{\alpha, q}(f)\right)(a, .)\right](b) . \tag{61}
\end{equation*}
$$

Proof. 1) Let $f$ be in $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right), a \in \mathbb{R}_{q}$ and $b \in \widetilde{\mathbb{R}}_{q}$. Using Corollary 3, we obtain $\Psi_{q, g}^{\alpha}(f)(a, b)=\sqrt{|a|} f * \overline{g_{a}}(b)=\sqrt{|a|} V_{\alpha, q}\left[{ }^{t} V_{\alpha, q}(f) *_{q} V_{\alpha, q}^{-1}\left(\overline{g_{a}}\right)\right](b)$. So, Theorem 7, Proposition 17 and the relation (55) achieve the proof.
2) Follows from Corollary 3, Theorem 8, and Propositions 14 and 17.

Theorem 11. Let $g$ be a q-wavelet, associated with the $q$-Dunkl operator, in $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$. Then for all $x \in \mathbb{R}_{q}$,

1) For all $f$ in $\mathcal{S}_{q,-1 / 2}\left(\mathbb{R}_{q}\right)$, we have

$$
\begin{aligned}
&{ }^{t} V_{\alpha, q}^{-1}(f)(x) \\
&= \frac{c_{\alpha, q}}{C_{\alpha, g}} \int_{-\infty}^{\infty}\left(\int _ { - \infty } ^ { \infty } V _ { \alpha , q } \left[\Phi_{q, K_{\alpha, q, 1} \circ}{ }^{t} V_{\alpha, q}(g)\right.\right. \\
&\left.(f)(a, .)](b) \times g_{(a, b), \alpha}(-x) \frac{|b|^{2 \alpha+1}}{|a|^{2 \alpha+3}} d_{q} b\right) d_{q} a ;
\end{aligned}
$$

2) For all $f$ in $\mathcal{S}_{q, \alpha}\left(\mathbb{R}_{q}\right)$, we have
$V_{\alpha, q}^{-1}(f)(x)=\frac{c_{-\frac{1}{2}, q}}{C_{g}} \int_{-\infty}^{\infty}\left(\int_{-\infty}^{\infty}{ }^{t} V_{\alpha, q}\left[\Psi_{q, K_{\alpha, q, 2}(g)}^{\alpha}(f)(a,).\right](b) g_{a, b}(-x) \frac{d_{q} b}{|a|^{2 \alpha+3}}\right) d_{q} a$.
Proof. The result derives from the previous theorem, Theorem 4 and ([4], Theorem 5 ).

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[^0]:    *Corresponding author. E-mail: neji.bettaibi@yahoo.fr
    †E-mail: rym.bettaieb@yahoo.fr
    ${ }^{\ddagger}$ E-mail: slimbouaziz@yahoo.fr

